



Recent results on branching Brownian motion on the positive real axis

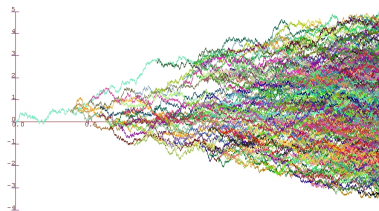
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Outline

- 1 Introduction
- 2 BBM with absorption
- 3 BBM with absorption, near-critical drift
- 4 BBM with absorption, critical drift

Branching Brownian motion (BBM)



Picture by Matt Roberts

Definition

- A particle performs **standard Brownian motion** started at a point $x \in \mathbb{R}$.
- With rate $1/2$, it **branches into 2 offspring** (can be generalized)
- Each offspring repeats this process independently of the others.

→ A **Brownian motion** indexed by a **tree**.

Why BBM ?

- Discrete counterpart: **branching random walk**, has lots of applications in diverse domains
 - Generalisation of age-dependent branching processes (*Crump-Mode-Jagers process*), model for asexual population undergoing mutation (position = fitness)
 - Toy model for *log-correlated field*, e.g. 2-dimensional Gaussian free field appearing notably in *Liouville quantum gravity theory*.
 - Used to study random walk in random environment on trees [Hu-Shi et al.](#), growth-fragmentation processes [Bertoin-Budd-Curien-Kortchemski](#), loop $O(n)$ model on random quadrangulations [Chen-Curien-M.](#), ...
- Intimate relation with (F-)KPP equation
- With diffusion constant depending on time : also known as [Derrida's](#) CREM spin glass model

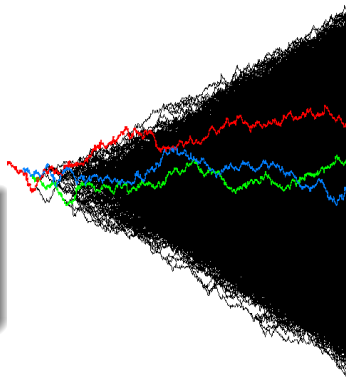
Maximum : LLN

M_t = maximum at time t .

LLN (Biggins '77)

Almost surely,

$$M_t/t \rightarrow 1, \quad \text{as } t \rightarrow \infty.$$



Picture by [Éric Brunet](#)

A family of martingales

For every $\theta \in \mathbb{R}$,

$$\mathbb{E}[\#\{u \in \mathcal{N}_t : X_u(t) \approx \theta t\}] = e^{\frac{1}{2}t} \mathbb{P}(B_t \approx \theta t) \approx e^{\frac{1}{2}(1-\theta^2)t}.$$

Martingales:

$$W_t^{(\theta)} = \sum_{u \in \mathcal{N}_t} e^{\theta X_u(t) - \frac{1}{2}(1+\theta^2)t}$$

Theorem (Biggins 78)

The martingale $(W_t^{(\theta)})_{t \geq 0}$ is *uniformly integrable* if and only if $|\theta| < 1$. In this case, for every $a, b \in \mathbb{R}$, $a < b$,

$$\frac{\#\{u \in \mathcal{N}_t : X_u(t) \in \theta t + [a, b]\}}{\mathbb{E}[\#\{u \in \mathcal{N}_t : X_u(t) \in \theta t + [a, b]\}]} \rightarrow W^{(\theta)} := W_\infty^{(\theta)}, \quad \text{a.s. as } t \rightarrow \infty.$$

Derivative martingale

For $\theta = 1$, $W_t^{(1)} \rightarrow 0$, almost surely as $t \rightarrow \infty$. **Derivative martingale:**

$$D_t = -\frac{d}{d\theta} W_t^{(\theta)} \Big|_{\theta=1} = \sum_{u \in \mathcal{N}_t} (t - X_u(t)) e^{X_u(t) - t}.$$

Theorem (Lalley-Sellke 87)

Almost surely, D_t converges as $t \rightarrow \infty$ to a non-degenerate r.v. D .

Theorem (Bramson 83 + Lalley-Sellke 87, Aïdekon 11)

Let M_t = maximum at time t . Then, conditioned on D , for some constant $C > 0$,

$$M_t - \left(t - \frac{3}{2} \log t\right) \Rightarrow \log CD + G,$$

where G is a standard Gumbel-distributed random variable.

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Absorption at the origin

- Start with one particle at $x \geq 0$.
- Add *drift* $-\mu$, $\mu \in \mathbb{R}$ to motion of particles.
- **Kill** particles upon hitting the origin.

Theorem (Kesten 78)

$$\mathbb{P}(\textit{survival}) > 0 \iff \mu < 1.$$

Why should we do this?

- Useful for the study of BBM **without** absorption (e.g., convergence of derivative martingale)
- **Biological interpretation**: natural selection
- Appears in other mathematical models, e.g. infinite bin models **Aldous**, **Mallein-Ramassany**

Absorption at the origin, $\mu \geq 1$

Start with one particle at 0, absorb particles at $-x$. N_x = number of particles absorbed at $-x$. Set

$$\theta_{\pm} = \mu \pm \sqrt{\mu^2 - 1}.$$

Theorem (Neveu 87, Chauvin 88)

$(N_x)_{x \geq 0}$ is a continuous-time Galton–Watson process. Moreover, almost surely as $x \rightarrow \infty$,

- If $\mu > 1$, $e^{-\theta_- x} N_x \rightarrow W^{(\theta_-)}$.
- If $\mu = 1$, $x e^{-x} N_x \rightarrow D$.

Theorem

As $x \rightarrow \infty$,

- $\mu > 1$: $\mathbb{P}(W^{(\theta_-)} > x) \sim C(\mu) x^{-\theta_+/\theta_-}$ Guivarc'h 90, Liu 00
- $\mu = 1$: $\mathbb{P}(D > x) \sim 1/x$ Buraczewski 09, Berestycki–Berestycki–Schweinsberg 10, M. 12

Absorption at the origin, $\mu \geq 1$ (contd.)

$$\theta_{\pm} = \mu \pm \sqrt{\mu^2 - 1}.$$

Theorem

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- $\mu = 1$: $\mathbb{P}(D > x) \sim 1/x$ *Buraczewski 09, Berestycki-Berestycki-Schweinsberg 10, M. 12*

Theorem (M. 10, Aïdekon-Hu-Zindy 12)

As $n \rightarrow \infty$,

- $\mu > 1$: $\mathbb{P}(N_x > n) \sim C(e^{\theta_+x} - e^{\theta_-x})/n^{-\theta_+/\theta_-}$.
- $\mu = 1$: $\mathbb{P}(N_x > n) \sim xe^x/(n(\log n)^2)$.

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Absorption at the origin, $\mu = 1 - \varepsilon$

Few works on $\mu < 1$ (Berestycki–Brunet–Harris–Miloś, Corre). But near-critical case $\mu = 1 - \varepsilon$, $0 < \varepsilon \ll 1$ well understood. Parametrize ε by

$$\varepsilon = \frac{\pi^2}{2L^2} \quad (\varepsilon \rightarrow 0 \iff L \rightarrow \infty).$$

Theorem (Brunet–Derrida 06, Gantert–Hu–Shi 08)

$$\mathbb{P}_1(\textit{survival}) = \exp(-(1 + o(1))L), \quad L \rightarrow \infty.$$

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Theorem (BBS 10)

There exists $C > 0$, such that, as $L \rightarrow \infty$,

$$\mathbb{P}_{L+x}(\text{survival}) \rightarrow 1 - \phi(x), \quad \phi(x) := \mathbb{E}[\exp(-CDe^x)].$$

and if $x = x(L)$ such that $L - x \rightarrow \infty$,

$$\mathbb{P}_x(\text{survival}) \sim C(L/\pi) \sin(\pi x/L) e^{x-L}.$$

Define

$$Z_t^L = \sum_{u \in \mathcal{N}_t} L \sin(\pi X_u(t)/L) e^{x-L}.$$

Then $(Z_t^L)_{t \geq 0}$ is (almost) a **martingale** for BBM with absorption at 0 and at L .

Theorem (BBS 10)

Suppose the initial configurations are such that $Z_0^L \Rightarrow z_0$ as $L \rightarrow \infty$, and $L - \max_u X_u(0) \rightarrow \infty$. Then $(Z_{L^{-3}t}^L)_{t \geq 0}$ converges as $L \rightarrow \infty$ (wrt fdis) to a continuous-state branching process started at z_0 . Moreover, $\mathbb{P}(\text{BBM survives forever}) \rightarrow \mathbb{P}(\text{CSBP started from } z_0 \text{ goes to } \infty)$.

The CSBP in the above theorem is Neveu's CSBP and has branching mechanism

$$\psi(u) = au + \pi^2 u \log u = a'u + \pi^2 \int_0^\infty (e^{-ux} - 1 + ux 1_{x \leq 1}) \frac{dx}{x^2},$$

for some (implicit) constants $a, a' \in \mathbb{R}$. In particular, it is supercritical (with ∞ mean).

BBS 10 proof (2)

Theorem (BBS 10)

If $x = x(L)$ such that $L - x \rightarrow \infty$,

$$\mathbb{P}_x(\text{survival}) \sim \frac{CL}{\pi} \sin(\pi x/L) e^{x-L}.$$

Proof: Set $w(x) := L \sin(\pi x/L) e^{x-L}$. Start BBM with $1/w(x)$ particles at x at time 0. Then

$$\mathbb{P}(\text{survival}) \rightarrow \mathbb{P}(\text{CSBP started at 1 goes to } \infty) \in (0,1).$$

Also, by independence,

$$1 - \mathbb{P}(\text{survival}) = (1 - \mathbb{P}_x(\text{survival}))^{1/w(x)} \sim \exp\left(-\frac{\mathbb{P}_x(\text{survival})}{w(x)}\right),$$

and so

$$\mathbb{P}_x(\text{survival}) \sim Cw(x). \quad \square$$

BBS 10 proof (3)

Theorem (BBS 10)

There exists $C > 0$, such that, as $L \rightarrow \infty$,

$$\mathbb{P}_{L+x}(\text{survival}) \rightarrow 1 - \phi(x), \quad \phi(x) = \mathbb{E}[\exp(-CDe^x)].$$

Proof: Wait a long time T (independent of L), so that $L - \max_u X_u(T) \gg 1$. Then using $L \sin(\pi x/L) \sim \pi(L-x)$ for $L-x \ll L$, we get

$$Z_T^L \approx \pi e^x D_T,$$

with $(D_t)_{t \geq 0}$ the derivative martingale of usual BBM. Let first $L \rightarrow \infty$ then $T \rightarrow \infty$ to get

$$\begin{aligned} \mathbb{P}_{L+x}(\text{survival}) &= 1 - \mathbb{E}[\mathbb{P}_{L+x}(\text{extinction} \mid \mathcal{F}_T)] \\ &\approx 1 - \mathbb{E}[\mathbb{P}(\text{CSBP started from } \pi e^x D_T \text{ goes to } 0)] \\ &\approx 1 - \mathbb{E}[\exp(-CDe^x)] = 1 - \phi(x). \quad \square \end{aligned}$$

BBS 10 convergence to CSBP

Basic idea

Decompose process into **bulk** + **fluctuations** by putting an additional absorbing barrier at L .

- **bulk**: Particles that don't hit L .
- **fluctuations**: Particles from the moment they hit L .

Then,

- $Z_t^{L,\text{bulk}}$ stays bounded over time scale L^3 .
- $Z_t^{L,\text{fluctuations}}$ increases from the contributions of the particles hitting L , an increase being roughly distributed as πD , with D derivative martingale limit.
- Particles hit L with rate $O(L^{-3})$.

BBS 10 convergence to CSBP

Basic idea

Decompose process into **bulk** + **fluctuations** by putting an additional absorbing barrier at $L - A$, where A is a large constant.

- **bulk**: Particles that don't hit $L - A$.
- **fluctuations**: Particles from the moment they hit $L - A$.

Then,

- $Z_t^{L,\text{bulk}}$ decreases almost deterministically as $\exp(-At/L^3)$.
- $Z_t^{L,\text{fluctuations}}$ increases from the contributions of the particles hitting L , an increase being roughly distributed as πe^{-AD} , with D derivative martingale limit.
- Particles hit $L - A$ with rate $O(e^A/L^3)$.

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- Particles hit $L - A$ with rate $O(e^A/L^3)$.

Recall: $\mathbb{P}(D > x) \sim 1/x$, $x \rightarrow \infty$. This yields convergence of $(Z_{L^3 t}^L)_{t \geq 0}$ to Neveu's CSBP as $L \rightarrow \infty$.

- 1 The basic phenomenological picture of BBM with near-critical drift (bulk + fluctuations) was established in [Brunet-Derrida-Mueller-Munier 06](#)
- 2 The techniques in [BBS 10](#) were a key ingredient in the study of BBM with selection of the N right-most particles, $N \gg 1$ ([M 16](#)). Relation between parameters: $\log N \approx L$, so $\varepsilon \approx \pi^2/2(\log N)^2$.

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Critical drift $\mu = 1$. Questions

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- Asymptotic of $\mathbb{P}_x(\text{survival until time } t)$?
- Conditioned on survival until time t , what does the BBM look like?

Critical drift $\mu = 1$. Questions

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- Conditioned on survival until time t , what does the BBM look like?

Kesten 78:

- Let $L_t = ct^{1/3}$, $c = (3\pi^2/2)^{1/3}$, Fix $x \geq 0$.

$$\mathbb{P}_x(\text{survival until time } t) = xe^{x-L_t+O((\log t)^2)}.$$

- Conditioned on survival until time t , with high probability,

$$\#\mathcal{N}_t \leq e^{O(t^{2/9}(\log t)^{2/3})} \quad \text{and} \quad \max_u X_u(t) \leq O(t^{2/9}(\log t)^{2/3}).$$

Note: $t^{1/3}$ scaling reminiscent of results about particles in BBM staying always close to the maximum [Faraud-Hu-Shi](#), [Fang-Zeitouni](#), [Roberts](#).

BBS 12 results

$$L_t = ct^{1/3}, \quad c = (3\pi^2/2)^{1/3}, \quad w_t(x) = L_t \sin(\pi x/L_t) e^{x-L_t}.$$

Theorem (BBS 12)

$$C_1 \leq \mathbb{P}_{L_t}(\text{survival until time } t) \leq C_2.$$

If $L_t - x \geq 1$,

$$C_1 w_t(x) \leq \mathbb{P}_x(\text{survival until time } t) \leq C_2 w_t(x).$$

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Theorem (Berestycki-M.-Schweinsberg, in preparation)

There exists $C > 0$, such that, as $t \rightarrow \infty$,

$$\mathbb{P}_{L_t+x}(\text{survival until time } t) \rightarrow 1 - \phi(x), \quad \phi(x) = \mathbb{E}[\exp(-CDe^x)].$$

and if $x = x(t)$ such that $L_t - x \rightarrow \infty$,

$$\mathbb{P}_x(\text{survival until time } t) \sim (C/\pi) w_t(x)$$

New results

$L_t = ct^{1/3}$, $c = (3\pi^2/2)^{1/3}$, $\zeta =$ time of extinction.

Corollary (BMS)

- 1 For fixed $x \in \mathbb{R}$, under \mathbb{P}_{L_t+x} , the r.v. $(\zeta - t)/t^{2/3}$ converges in law to $\frac{3}{c}(G - x - \log CD)$, where G is a Gumbel-distributed random variable independent of D .
- 2 Suppose $L_t - x \rightarrow \infty$. Conditionally on $\zeta > t$, under \mathbb{P}_x , $(\zeta - t)/t^{2/3}$ converges in law to $\text{Exp}(c/3)$ as $t \rightarrow \infty$.

Reason: For fixed $s \geq 0$,

$$L_{t+st^{2/3}} = L_t + \frac{c}{3}s + o(1).$$

This gives as $t \rightarrow \infty$, for fixed $x \in \mathbb{R}$,

$$\mathbb{P}_{L_t+x}(\zeta \leq t + st^{2/3}) \rightarrow \phi(x - \frac{c}{3}s) = \mathbb{E}[e^{-CDe^{x-(c/3)s}}].$$

New results (contd.)

$L_t = ct^{1/3}$, $c = (3\pi^2/2)^{1/3}$, ζ = time of extinction, $M_t = \max_u X_u(t)$.

Theorem (BMS)

- 1 For fixed $x \in \mathbb{R}$, under \mathbb{P}_{L_t+x} , the r.v. $M_t/t^{2/9}$ converges in law to $(3c^2(G - x - \log CD) \vee 0)^{1/3}$, where G is a Gumbel-distributed random variable independent of D .
- 2 Suppose $L_t - x \rightarrow \infty$. Conditionally on $\zeta > t$, under \mathbb{P}_x , $M_t/t^{2/9}$ converges in law to $(3c^2V)^{1/3}$, where $V \sim \text{Exp}(1)$.

Reason: morally, $M_t \approx L_{\zeta-t}$ if $\zeta > t$ (and $M_t = 0$ if $\zeta \leq t$).

New results (contd.)

$L_t = ct^{1/3}$, $c = (3\pi^2/2)^{1/3}$, ζ = time of extinction, $M_t = \max_u X_u(t)$.

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Reason: morally, $M_t \approx L_{\zeta-t}$ if $\zeta > t$ (and $M_t = 0$ if $\zeta \leq t$).

Same result holds with M_t replaced by $\log \#\mathcal{N}_t$.

New results (contd.)

- $L_t(s) = L_{t-s} = c(t-s)^{1/3}$.
- $Z_t(s) = \sum_{u \in \mathbb{N}_s} w_{t-s}(X_u(s)) = \sum_{u \in \mathbb{N}_s} L_t(s) \sin(\pi X_u(s)/L_t(s)) e^{X_u(s) - L_t(s)}$.

Theorem (BMS)

Suppose the initial configurations are such that $Z_t(0) \Rightarrow z_0$ as $t \rightarrow \infty$, and $L_t - \max_u X_u(0) \rightarrow \infty$. Then

- $(Z_t(t(1 - e^{-s})))_{s \geq 0}$ converges as $t \rightarrow \infty$ (wrt fdis) to the CSBP with branching mechanism $\psi(u) = au + \frac{2}{3}u \log u$ started at z_0 .
- $\mathbb{P}(\zeta > t) \rightarrow \mathbb{P}(\text{CSBP started from } z_0 \text{ goes to } \infty)$, as $t \rightarrow \infty$.
- Conditioned on $\zeta > t$, $(Z_t(t(1 - e^{-s})))_{s \geq 0}$ converges as $t \rightarrow \infty$ to the CSBP started at z_0 conditioned to go to ∞ .

New results (contd.)

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- $\mathbb{P}(\zeta > t) \rightarrow \mathbb{P}(\text{CSBP started from } z_0 \text{ goes to } \infty)$, as $t \rightarrow \infty$.
- Conditioned on $\zeta > t$, $(Z_t(t(1 - e^{-s})))_{s \geq 0}$ converges as $t \rightarrow \infty$ to the CSBP started at z_0 conditioned to go to ∞ .

Proof inspired by [BBS 10](#) but requiring furthermore precise estimates for density of Brownian motion in curved domains refining those obtained in [Roberts 12](#).

Relation between results

In order to understand the relation between the several results, we use the **long-time behavior** of Neveu's CSBP. It grows doubly-exponentially:

Theorem (Neveu 92)

Let $(Y_t)_{t \geq 0}$ be the CSBP with branching mechanism $\psi(u) = au + bu \log u$, $a \in \mathbb{R}$, $b > 0$, starting at $z_0 > 0$. Then,

$\frac{\log Y_t}{e^{bt}}$ converges almost surely to a limit Y .

In particular, almost surely, the process **survives** iff $Y > 0$. Furthermore, there is $C = C(a, b)$, such that $Y - \log Cz_0$ follows the **Gumbel** distribution.

Relation between results (contd.)

Heuristic: As long as $R_s \approx L_t(s)$, we expect $\log Z_t(s) \approx R_s - L_t(s)$. When does R_s become significantly different from $L_t(s)$?

Relation between results (contd.)

Heuristic: As long as $R_s \approx L_t(s)$, we expect $\log Z_t(s) \approx R_s - L_t(s)$. When does R_s become significantly different from $L_t(s)$?

Answer: With the asymptotic growth of Neveu's CSBP, can check that $\log Z_t(s) \ll L_t(s)$ as long as $t - s \gg t^{2/3}$, hence the turning point is at $s = t - Kt^{2/3}$ for K large and one can read off M_t as well as $(\zeta - t)/t^{2/3}$ from $Z_t(s)$ at that point.

Conclusion

- 1 We were able to push the techniques from [BBS 10](#) on BBM with near-critical drift to the case of critical drift.
- 2 Results might be of help for the fine study of other models involving extremal particles of BBM.

Example [CREM](#) ([Derrida's](#) continuous random energy model): BBM during time $[0, T]$ with time-dependent diffusion constant $2\sigma^2(t/T)$. If σ^2 is strictly decreasing, then ([M., Zeitouni 16](#)) there exists a function $m(T)$ and constants $c, c', c'' > 0$, such that

$$\{\text{maximum at time } T\} - m(T) \Rightarrow \text{mixture of Gumbel,}$$
$$\text{with } m(T) = cT - c'T^{2/3} - c'' \log T + O(1).$$

Removing the $O(1)$ term would require an analysis similar to the one performed here.

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- 2 Results might be of help for the fine study of other models involving extremal particles of BBM.

Example CREM (Derrida's continuous random energy model): BBM during time $[0, T]$ with time-dependent diffusion constant $2\sigma^2(t/T)$. If σ^2 is strictly decreasing, then (M., Zefouni 16) there exists a function $m(T)$ and constants $c, c', c'' > 0$, such that

{maximum at time T } $- m(T) \Rightarrow$ mixture of Gumbel,

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Removing the $O(1)$ term would require an analysis similar to the one performed here.